

# Effect of Substrate on the Pre-Steady-State Kinetics of the Na<sup>+</sup>/Glucose Cotransporter

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**ABSTRACT** When measuring Na<sup>+</sup>/glucose cotransporter (SGLT1) activity in *Xenopus* oocytes with the two-electrode voltage-clamp technique, pre-steady-state currents dissipate completely in the presence of saturating  $\alpha$ -methyl-glucose ( $\alpha$ MG, a nonhydrolyzable glucose analog) concentrations. In sharp contrast, two SGLT1 mutants (C255A and C511A) that lack a recently identified disulfide bridge express the pre-steady-state currents in the presence of  $\alpha$ MG. The dose-dependent effects of  $\alpha$ MG on pre-steady-state currents were studied for wild-type (wt) SGLT1 and for the two mutants. Increases in  $\alpha$ MG concentration reduced the total transferred charge (partially for the mutants, totally for wt SGLT1), shifted the transferred charge versus membrane potential (Q-V) curve toward positive potentials, and significantly modified the time constants of the pre-steady-state currents. A five-state kinetic model is proposed to quantitatively explain the effect of  $\alpha$ MG on pre-steady-state currents. This analysis reveals that the reorientation of free transporter is the slowest step for wt SGLT1 either in the presence or in the absence of  $\alpha$ MG. In contrast, the conformational change of the fully loaded mutant transporters constitutes their rate-limiting step in the presence of substrate and explains the persistence of pre-steady-state currents in this situation.

## INTRODUCTION

The Na<sup>+</sup>/glucose cotransporter SGLT1 is a member of the SLC5 family and has been the archetype of this class of Na<sup>+</sup>-coupled substrate transporters. Soon after its cloning (1), expression in *Xenopus* oocytes enabled the measurement of pre-steady-state currents, i.e., transient currents observed in the absence of substrate, which were suggestive of gating currents observed in voltage-dependent channels. As these currents were Na<sup>+</sup> dependent and were absent in the presence of the specific inhibitor phlorizin (Pz) or in the presence of glucose, they were considered to represent charge displacements occurring during the voltage-dependent reorientation of the Na<sup>+</sup>-binding site and upon Na<sup>+</sup> binding (2).

Pre-steady-state currents have been extremely useful for devising a credible transport mechanism, with quantitative estimation of the rate constants linking the different conformational states. The original model (2) proposed in 1992 has been challenged both theoretically and experimentally ((3–8); for review see Wright and Turk (9)), and new steps have been proposed. Recently Loo et al., using fluorescently labeled mutants (10), have reported extremely slow conformational changes (time constants on the order of 100 ms) which have yet to be quantitatively explained by any proposed kinetic model. In particular, this observation is incompatible with a rate-limiting step of 50 s<sup>-1</sup>, which was proposed for the translocation of the fully loaded transporter in the human isoform of SGLT1 (hSGLT1) (2,11).

One characteristic observed for nearly all cotransporters studied (Na<sup>+</sup>/glucose cotransporters SGLT1 and SGLT2, Na<sup>+</sup>/myo-inositol cotransporters SMIT1 and SMIT2, Na<sup>+</sup>/monocarboxylate cotransporter SMCT1, Na<sup>+</sup>/Pi cotransporter NaPiII, Na<sup>+</sup>/I<sup>-</sup> symporter NIS, gamma-aminobutyric acid transporters GAT1 and GAT3, H<sup>+</sup>/hexose cotransporter STP1, and Cl<sup>-</sup>-dependant K<sup>+</sup>/amino acids transporter KAAT1) is that addition of a saturating concentration of substrate leads to the total inhibition of pre-steady-state currents (11–25). Surprisingly, with the exception of GAT1, no quantitative explanation has been proposed to explain this phenomenon.

Recently, we identified a disulfide bridge between C255 and C511 in hSGLT1 (26). An interesting feature of mutants C255A and C511A, which we have not previously published, is that they express pre-steady-state currents in the presence of a saturating  $\alpha$ MG concentration, in contrast to what is observed with wild-type (wt) SGLT1. This phenomenon has prompted us to examine the dose-dependent effects of  $\alpha$ MG on the pre-steady-state currents for the two mutants as well as for wt SGLT1 and to propose a quantitative explanation using a kinetic model displaying different rate-limiting steps for the wt SGLT1 and the mutant cotransporters.

## MATERIALS AND METHODS

### Oocyte preparation and injection

Oocytes were surgically removed from *Xenopus laevis* frogs, dissected, and defolliculated as described previously (27). One day after defolliculation, oocytes were injected with 46 nl of water containing mRNA (0.1  $\mu$ g/ $\mu$ l and 0.25  $\mu$ g/ $\mu$ l for wt SGLT1 and mutants, respectively) to obtain maximal protein expression. Oocytes were maintained in Barth's solution (in mM: 90 NaCl, 3 KCl, 0.82 MgSO<sub>4</sub>, 0.41 CaCl<sub>2</sub>, 0.33 Ca(NO<sub>3</sub>)<sub>2</sub>, 5 Hepes, pH 7.6) supplemented with 5% horse serum, 2.5 mM Na<sup>+</sup> pyruvate, 100 units/ml

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penicillin, and 0.1 mg/ml streptomycin for 4–7 days before electrophysiological experimentation.

## Molecular biology

The constructions prepared for obtaining the mutants C255A and C511A have been described elsewhere (26).

## Electrophysiology

The saline solution normally used in our electrophysiological experiments is composed of (in mM): 90 NaCl, 3 KCl, 0.82 MgCl<sub>2</sub>, 0.74 CaCl<sub>2</sub>, and 10 Hepes and the pH was adjusted to 7.6 with NaOH. Two-electrode voltage-clamp experiments were performed using an Oocyte Clamp OC-725 (Warner Instruments, Hamden, CT) and a data acquisition system (Digidata 1322A and Clampex 8.2, Axon Instruments, Union City, CA). Current and voltage microelectrodes were filled with 1 M KCl and had a resistance of 1–2 M $\Omega$ . The bath current electrode was an Ag–AgCl pellet, and the reference electrode was a 1 M KCl agar bridge. The oocytes were clamped to a membrane potential ( $V_m$ ) of  $-50$  mV, and three repetitions of  $V_m$  steps between  $+70$  and  $-170$  mV (by increments of 20 mV, 300 ms duration, no series resistance compensation used) were applied with an interval of 1.7 s between each step. Ninety-five percent of the command voltage step was reached in 3–4 ms. Data were obtained with a sampling frequency of 10 kHz, without filtering, and the three repetitions were averaged for each experiment.

## Data analysis

Pre-steady-state current analysis was performed as described previously (26). Briefly, the transferred charge was obtained at each membrane potential (Q-V curve) by subtracting the integrated baseline-corrected currents in Pz solution (200  $\mu$ M) from similar currents in saline solution (either in the presence or absence of  $\alpha$ MG). Thus, the transferred charge calculated corresponds to the total charge in one experimental condition minus the total charge in the presence of Pz. The total charge in the presence of Pz was found to be linear with voltage as expected if it was mainly due to the presence of the oocyte capacitive current. The baseline correction was obtained from the mean current measured between 50 and 80 ms after the initiation of a voltage pulse. A simple Boltzmann equation was fitted to the Q-V curve to estimate  $V_{1/2}$  (the voltage at which half of the charge is transferred),  $Q_{\max}$  (the amplitude of the total charge transferred), and  $z$  (the valence of the transferred charge) (26). The time constants ( $\tau_{\text{slow}}$ ) were evaluated by fitting a double exponential on the  $I_{\text{transit}}$  ( $I_{\text{saline}} - I_{\text{Pz}}$ ) with the Clampfit 8.2 program (Axon Instruments). Only the slow time constant ( $\tau_{\text{slow}}$ ; 2–10 ms), which has the dominant amplitude, was considered.

## Statistics

Experiments were performed on at least six oocytes obtained from a minimum of two different donors. Data are reported as means  $\pm$  SE and are compared using unpaired Student's *t*-test; statistical significance was set at  $P < 0.05$ . Errors bars were omitted when smaller than the symbol size.

## RESULTS

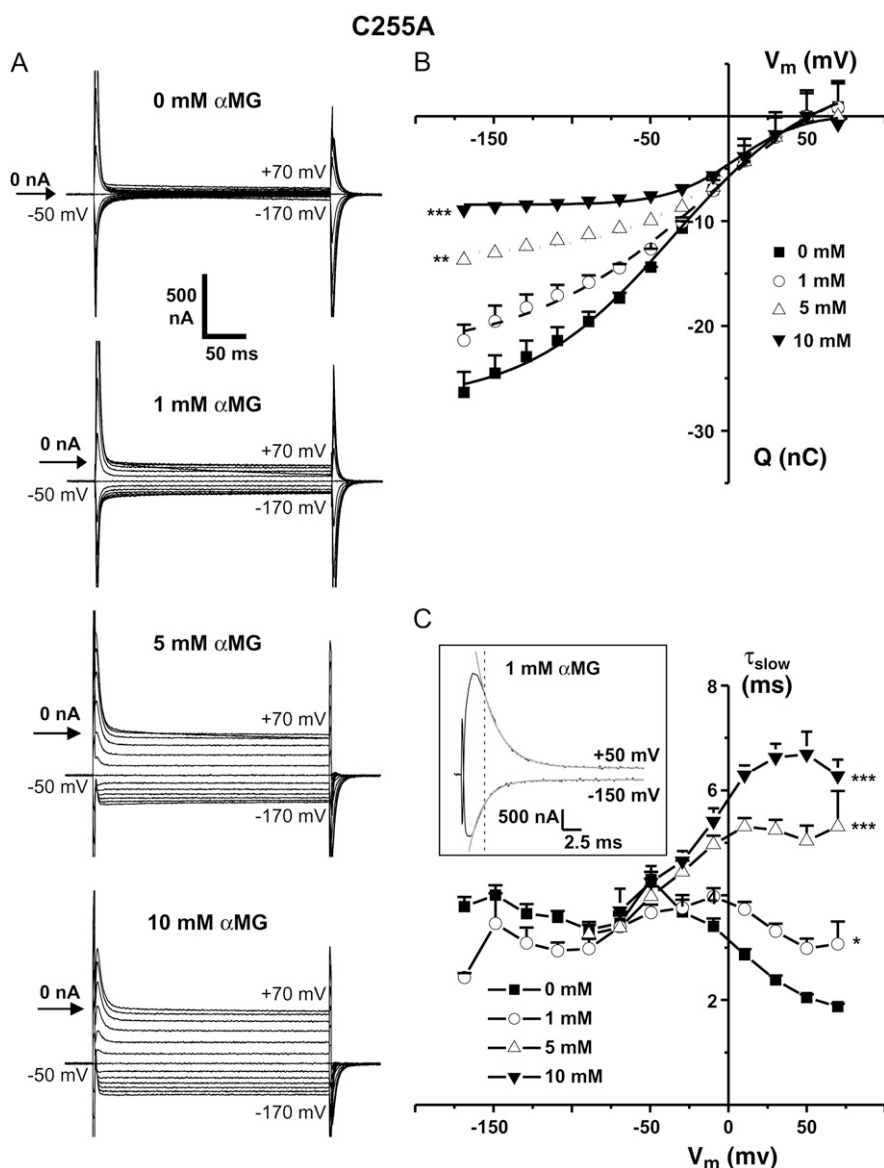
### Pre-steady-state currents in the presence of $\alpha$ MG for C255A and C511A

In contrast to wt SGLT1, the mutants C255A and C511A clearly exhibit pre-steady-state currents in the presence of a saturating  $\alpha$ -methyl-glucose ( $\alpha$ MG) concentration (26), par-

ticularly at depolarizing  $V_m$  levels. Figs. 1 A and 2 A show the Pz-sensitive currents with different voltage steps for the two mutant proteins using several  $\alpha$ MG concentrations (0, 1, 5, and 10 mM  $\alpha$ MG). As the capacitive currents are eliminated by subtraction of currents measured in the presence of Pz, the transient currents at each voltage step directly represent SGLT1-specific pre-steady-state currents. The integrals of the transient currents measured at different  $V_m$  are used to produce transferred charge versus  $V_m$  (Q-V) curves. The Q-V curves for the two mutants are shown in Figs. 1 B and 2 B. As the cotransporter conformation at very positive  $V_m$  levels is predicted to be independent of the presence of extracellular  $\alpha$ MG (the binding site being predicted to face inside in all cases (2,4,10)), each Q-V curve was shifted vertically to have  $Q = 0$  at  $+50$  mV under all conditions. This allows direct comparison of the transferred charge at different  $V_m$ ; it is clear that they decrease as the  $\alpha$ MG concentration increases. It is also clear from Figs. 1 B and 2 B that the voltage range over which charge can be transferred is reduced in amplitude and displaced toward more positive potentials when the  $\alpha$ MG concentration is increased. The measured charge, at saturating  $\alpha$ MG concentration, has reached its plateau level at  $-50$  mV and remains basically constant as  $V_m$  becomes more negative. However, a saturating  $\alpha$ MG concentration does not totally abolish the transferred charge. A Boltzmann relation can be fitted to the Q-V curves, which yields a value for the parameter  $V_{1/2}$ , the  $V_m$  at which half of the mobile charge is equally balanced between the inward and outward facing positions. For both mutants, an increase in  $\alpha$ MG concentration shifts the  $V_{1/2}$  toward more positive values. For mutant C255A, the measured  $V_{1/2}$  averaged  $-35 \pm 5$  mV at 0 mM and  $5 \pm 5$  mV at 10 mM  $\alpha$ MG ( $n = 6$ ). For mutant C511A, the corresponding values were  $-28 \pm 2$  mV at 0 mM and  $13 \pm 3$  mV at 10 mM  $\alpha$ MG ( $n = 6$ ).

The progressive addition of extracellular  $\alpha$ MG also affected the  $\tau_{\text{slow}}$  for pre-steady-state currents as depicted in Figs. 1 C and 2 C. They were obtained by fitting double exponentials to the Pz-sensitive currents. As reported previously (26), in the absence of  $\alpha$ MG,  $\tau_{\text{slow}}$  for the mutant proteins reached a plateau of 4–5 ms at hyperpolarized  $V_m$ , which is about half of the value measured for the wt SGLT1. An increase in  $\alpha$ MG concentration clearly produced an increase of  $\tau_{\text{slow}}$  at positive  $V_m$ . For  $V_m$  more negative than  $-50$  mV, addition of  $\alpha$ MG accelerated the transient currents, as shown in Fig. 2 C for 1 mM  $\alpha$ MG. The currents in the presence of higher concentrations can no longer be fitted accurately at these voltages (see Fig. 2 C for 10 mM  $\alpha$ MG and Fig. 1 C for 5 and 10 mM  $\alpha$ MG). Consequently, at 1 mM  $\alpha$ MG, the  $\tau_{\text{slow}}$  versus  $V_m$  ( $\tau_{\text{slow}}-V$ ) curve has a bell shape centered around  $-50$  mV for both mutants. At 10 mM  $\alpha$ MG,  $\tau_{\text{slow}}$  starts at very low values for negative  $V_m$  and reaches a value of 6–7 ms for both mutants at depolarizing potentials.

The stability of the preparation as a function of time was tested in each experiment by comparing pre-steady-state currents in the absence of  $\alpha$ MG measured before and after



**FIGURE 1** Pre-steady-state currents of mutant C255A in the presence of different  $\alpha$ MG concentrations. (A) Pre-steady-state current traces at different  $V_m$  in the presence of various  $\alpha$ MG concentrations (0, 1, 5, and 10 mM) for a typical C255A-expressing oocyte. The currents were obtained by subtracting the currents in the presence of 200  $\mu$ M Pz from the currents measured in the various conditions. (B) Q-V curves in different  $\alpha$ MG concentrations were compared to those in the absence of substrate. Values were shifted to have the same  $Q = 0$  at  $V_m = +50$  mV. The curve represents a Boltzmann equation fitted to the points ( $n = 6$ ). (C) Time constants of the pre-steady-state currents in different  $\alpha$ MG concentrations were compared to those in the absence of substrate ( $n = 6$ ). The inset represents the double exponential fit (gray line) of the Pz-sensitive currents (black line) shown in panel A at 1 mM  $\alpha$ MG for the indicated  $V_m$ . The dotted line indicates the time at which 95% of the  $V_m$  is achieved. Means  $\pm$  SE are shown. Stars indicate statistical significance with respect to the values in 0 mM  $\alpha$ MG (\*,  $P \leq 0.05$ ; \*\*,  $P \leq 0.01$ ; \*\*\*,  $P \leq 0.001$ ).

having presented the different  $\alpha$ MG concentrations. In all cases, the Q-V and  $\tau_{\text{slow}}$ -V curves were found identical. In addition, the effects of  $\alpha$ MG on the pre-steady-state currents were found to be independent on the order of the  $\alpha$ MG concentrations applied (increasing or decreasing concentrations).

### Pre-steady-state currents in the presence of $\alpha$ MG for wt SGLT1

Although the inhibitory effect of  $\alpha$ MG on wt SGLT1 pre-steady-state currents has long been known (11,25), to our knowledge it has never been studied quantitatively in a dose-dependent manner nor has it been explained with the use of a kinetic model. Given our findings with the two SGLT1 mutant proteins, we sought to characterize this effect in detail on wt SGLT1 and to compare it with that observed for the mutants. Fig. 3 A shows Pz-sensitive pre-steady-state currents recorded

from wt SGLT1 in the absence or in the presence of different  $\alpha$ MG concentrations (0.5 mM, 1 mM, and 5 mM). In agreement with previous reports, addition of extracellular  $\alpha$ MG leads to a progressive decrease in the pre-steady-state currents and to the appearance of steady-state inward  $\text{Na}^+$ /glucose flux, which have been described in detail (11,25). The Q-V curves obtained with different  $\alpha$ MG concentrations is illustrated in Fig. 3 B. It is obvious that a saturating concentration of  $\alpha$ MG (5 mM) completely abolished charge transfer ( $n = 9$ ), at least within the time resolution provided by the two-electrode voltage-clamp technique. In the presence of 5 mM  $\alpha$ MG, the amplitude of the transferred charge that can be detected is approximately equal to that measured from a noninjected oocyte ( $<1$  nC). The time constants of these pre-steady-state currents are shown in Fig. 3 C. Increased  $\alpha$ MG concentrations produce a clear acceleration of the transient currents at negative  $V_m$ , whereas the value of the time constant remains

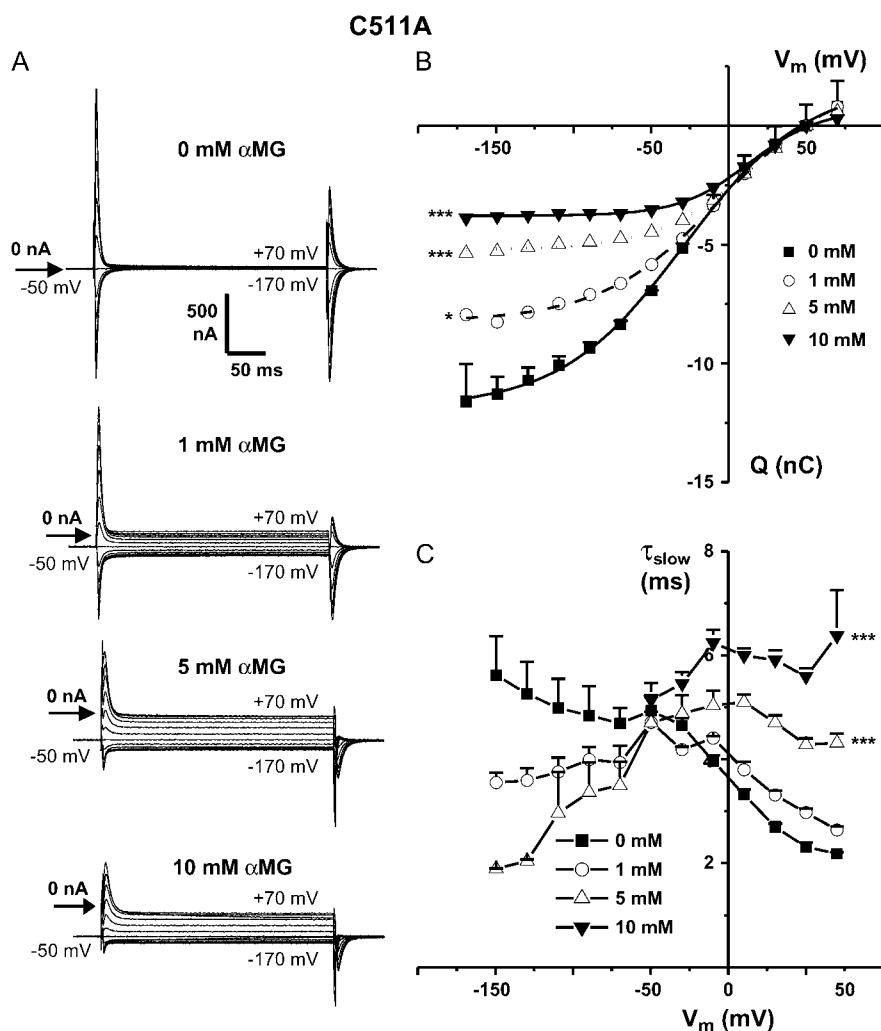


FIGURE 2 Pre-steady-state currents of mutant C511A in the presence of different  $\alpha$ MG concentrations. (A) Pre-steady-state current traces at different  $V_m$  in the presence of various  $\alpha$ MG concentrations (0, 1, 5, and 10 mM) for a typical C511A-expressing oocyte. The currents were obtained by subtracting the currents in the presence of 200  $\mu$ M Pz from the currents measured in the various conditions. (B) Q-V curves in different  $\alpha$ MG concentrations were compared to those in the absence of substrate. Values were shifted to have the same  $Q = 0$  at  $V_m = +50$  mV. The curve represents a Boltzmann equation fitted to the points ( $n = 6$ ). (C) Time constants of the pre-steady-state currents in different  $\alpha$ MG concentrations were compared to those in the absence of substrate. ( $n = 6$ ). Means  $\pm$  SE are shown. Stars indicate statistical significance (see Fig. 1 legend).

the same at positive  $V_m$ . In the presence of 5 mM  $\alpha$ MG, the transient currents are difficult to fit to a double exponential equation because  $\tau_{slow}$  shows only modest voltage dependence and reaches a value of  $\sim 2.5$  ms (at +70 mV) with very small amplitude ( $n = 7$ ). Both parameters are close to the limits inherent to the two-electrode voltage-clamp technique.

Fig. 4 summarizes the effects of  $\alpha$ MG on the  $V_{1/2}$  and on the normalized amplitude of the total transferred charge for wt SGLT1 versus the two mutants. In Fig. 4 A, the shift in  $V_{1/2}$  produced by the addition of  $\alpha$ MG to the wt SGLT1 is compared with the shifts mentioned above for the two mutants. For wt SGLT1,  $V_{1/2}$  goes rapidly from  $-60 \pm 5$  mV at 0 mM to  $12 \pm 12$  mV at 1 mM  $\alpha$ MG ( $n = 9$ ) and, at  $\alpha$ MG concentrations higher than 1 mM, the amplitude of the Q-V curve is reduced to such an extent that the fitting of a Boltzmann curve is not reliable. In all cases, the  $V_{1/2}$  is progressively shifted toward positive  $V_m$  levels as the external [ $\alpha$ MG] increases, but this shift is less marked for the mutants than for wt SGLT1.

To quantitatively compare the sensitivity of the pre-steady-state currents to  $\alpha$ MG, the decreases in transferred charge

caused by  $\alpha$ MG were normalized to the total transferred charge in the absence of substrate (i.e.,  $(Q_{saline} - Q_{\alpha MG})/Q_{saline}$ ). A simple Michaelis-Menten equation was fitted to  $(Q_{saline} - Q_{\alpha MG})/Q_{saline}$  taken at  $-170$  mV as a function of  $\alpha$ MG concentration after setting  $Q = 0$  for  $V_m = +50$  mV (as done in Figs. 1 B, 2 B, and 3 B). It is clear that for the wt SGLT1, a high  $\alpha$ MG concentration inhibits 100% of the transferred charge, whereas only partial inhibition ( $\sim 65$ – $75\%$  inhibition) was observed for the mutant proteins. It was found that the  $\alpha$ MG-sensitive charge is consistent with  $\alpha$ MG affinity constants ( $K_{mQ}^{\alpha MG}$ ) of  $0.48 \pm 0.05$  mM for wt SGLT1 (see Fig. 4 B),  $5 \pm 2$  mM for mutant C255A, and  $2 \pm 1$  mM for mutant C511A whereas their  $K_{ml}^{\alpha MG}$  constants, using  $\alpha$ MG-sensitive cotransport current (steady-state values,  $I_{ss}(\alpha MG)$ ), were  $0.97 \pm 0.1$  mM for wt SGLT1 (26,27),  $1.6 \pm 0.2$  mM (C255A), and  $1.6 \pm 0.2$  mM (C511A) at  $-170$  mV (26).

### Kinetic model for pre-steady-state currents

A scheme of the simple five-state kinetic model used is presented in Fig. 5. The substrate-binding and -debinding

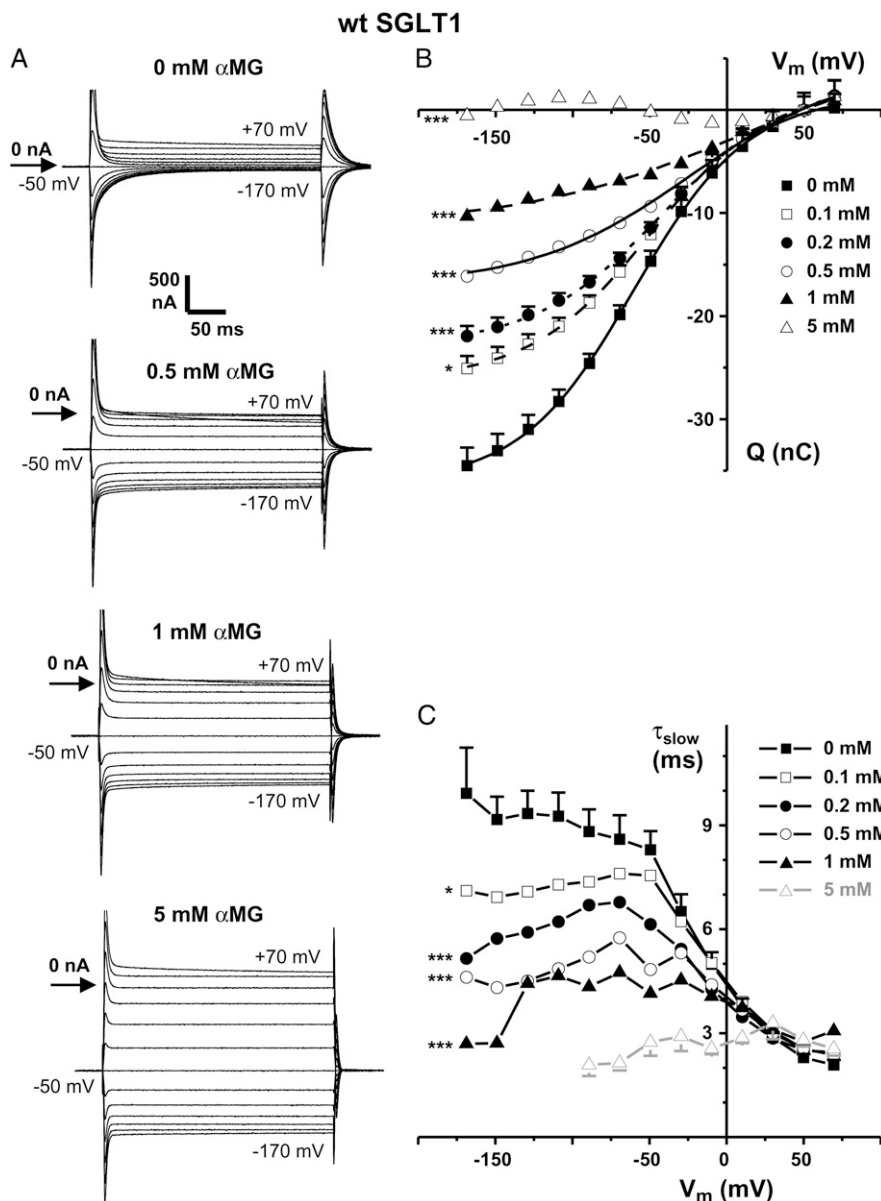


FIGURE 3 Pre-steady-state currents of wt SGLT1 in the presence of different αMG concentrations. (A) Pre-steady-state current traces for different  $V_m$  in the presence of various αMG concentrations (0, 0.5, 1, and 5 mM) containing solutions for a typical wt SGLT1-expressing oocyte. The currents were obtained by subtracting the currents in the presence of 200 μM Pz to the ones in the various conditions. (B) Q-V curves in different αMG concentrations as compared to those in the absence of substrate. Values were shifted to have the same  $Q = 0$  at  $V_m = +50$  mV. The curve represents a Boltzmann relation fitted to the points ( $n = 9$ ). (C) Time constants of the pre-steady-state currents in different αMG concentrations as compared to those in the absence of substrate. The time constants at 5 mM αMG were plotted in gray because of the small amplitude of the exponential giving uncertainty about the value of this time constant ( $n = 7$ ). Means  $\pm$  SE are shown. Stars indicate the statistical significance (see Fig. 1 legend).

steps ( $k_{45}$  and  $k_{54}$ ) are voltage independent (2,28) as are the lumped reactions  $k_{41}$  and  $k_{51}$ , which represent the conformational changes of the  $\text{Na}^+$ -loaded transporter (involved in the leak current) and the  $\text{Na}^+$ - and αMG-loaded transporter, respectively, with their associated intracellular release steps. It was found that the extracellular  $\text{Na}^+$ -binding reaction could be assumed to be a fast reaction at equilibrium without losing any fitting performance. The voltage-dependent reactions are expressed as follows (for  $i = 1, 2, 3; j = i + 1$ ):

$$k_{ij} = k_{ij0} \exp\left(z_i \alpha_i \frac{FV_m}{RT}\right), k_{ji} = k_{ji0} \exp\left(-z_i (1 - \alpha_i) \frac{FV_m}{RT}\right), \quad (1)$$

where  $z_i$  is the valence of the equivalent moving charge,  $\alpha_i$  represents the asymmetry of the energy barrier,  $V_m$  is the

membrane potential, and  $F$ ,  $R$ , and  $T$  have their usual meanings (see also Table 1).

The affinity for intracellular  $\text{Na}^+$  ( $K_m^{\text{Na}^+}$ ) was previously estimated for rabbit SGLT1 in giant, excised, inside-out patches and was found to vary from 44 to 70 mM (28,29). In agreement with these estimates, we recently found that intracellular  $\text{Na}^+$  has to be increased by blocking the  $\text{Na}^+/\text{K}^+$ -ATPase overnight to generate a measurable outward  $\text{Na}^+/\text{glucose}$  current upon intracellular glucose injection (30). This is consistent with an intracellular  $K_m^{\text{Na}^+}$  much higher than the physiological intracellular  $\text{Na}^+$  concentration. On the other hand, the estimation of  $K_m^{\alpha\text{MG}}$  for the intracellular site is  $\sim 35$  mM (28,29). Consequently, under physiological conditions, the inverse mode of transport is highly unlikely, which is reflected by very low values of  $k_{14}$  and  $k_{15}$  (see

Table 1). Thus, we assumed that the probability of finding the intracellular site loaded with  $\text{Na}^+$  and  $\alpha\text{MG}$  was negligible, and we reduced the potential seven-state kinetic model into the five-state kinetic model shown in Fig. 5.

The numerical simulations were performed using MATLAB 6.5.0 software (MathWorks, Natick, MA). Transferred charges were calculated as the integral of the pre-steady-state currents as done during the analysis of the experimental data and were also fitted with a simple Boltzmann curve to deduce the  $V_{1/2}$ ,  $Q_{\text{max}}$ , and  $z$  parameters. As the analytical expression for the time constant in the presence of substrate cannot be obtained in a five-state model, the numerical values of the time constant were obtained by taking the reciprocal of the eigenvalues of the following matrix, as previously reported (10):

$$\frac{d}{dt} \begin{bmatrix} C_1 \\ C_2 \\ C_3 \\ C_4 \\ C_5 \end{bmatrix} = \begin{bmatrix} -(k_{12} + k_{15}) & k_{21} & 0 & 0 & 0 \\ k_{12} & -(k_{21} + k_{23}) & k_{32} & 0 & 0 \\ 0 & k_{23} & -(k_{32} + k_{34}) & k_{43} & 0 \\ k_{14} & 0 & k_{34} & -(k_{45} + k_{43} + k_{41}) & k_{54} \\ k_{15} & 0 & 0 & k_{45} & -(k_{54} + k_{51}) \end{bmatrix} \begin{bmatrix} C_1 \\ C_2 \\ C_3 \\ C_4 \\ C_5 \end{bmatrix}.$$

Moreover, the model predicts the relaxation of pre-steady-state currents toward steady-state currents. The current ( $I$ ) versus  $V_m$  curves are sigmoid, as observed for the experimental current versus membrane potential ( $I$ - $V$ ) curves in the absence and in the presence of  $\alpha\text{MG}$  (not shown). The rate constants of the model were adjusted by trial and error to obtain a satisfactory fit to the measured  $Q$ - $V$  curves, the  $V_{1/2}$ , and the  $\tau_{\text{slow}}$ s of the pre-steady-state currents. The steady-state parameters ( $I$ - $V$  curves,  $K_m^{\text{Na}^+}$ , or  $K_m^{\alpha\text{MG}}$ ) were not considered as criteria for the adjustment of the model parameters but were found afterwards to be in accordance with the experimental values.

### Simulation of pre-steady-state currents in the absence of $\alpha\text{MG}$

Table 1 gives the rate constants used by Chen et al. (4) and by Loo et al. (11) as well as the rate constants used by this study to simulate the time constants of the pre-steady-state currents. As shown in Fig. 6 A, our new set of rate constants can reproduce, in a generally satisfying manner, the currents, the  $Q$ - $V$  curves, and the  $\tau_{\text{slow}}$ - $V$  curves for wt SGLT1 exposed to different  $\alpha\text{MG}$  concentrations. In the presence of 90 mM  $\text{Na}^+$ , a  $k_{210}/k_{120}$  ratio of  $\sim 0.5$  is required for the plateau effect seen on the  $\tau_{\text{slow}}$ - $V$  curve at hyperpolarizing  $V_m$ .  $k_{23}$  is a crucial rate constant because it is voltage independent and becomes rate limiting for  $V_m$  below  $-70$  mV (at 90 mM  $\text{Na}^+$ ). Indeed, the value of  $1/k_{23}$  closely corresponds to the plateau

value reached by  $\tau_{\text{slow}}$  at very negative  $V_m$ . The ratio  $k_{23}/k_{32}$  is also responsible for the voltage dependence of  $\tau_{\text{slow}}$  observed at depolarizing  $V_m$ . At 0 mV,  $K_{34}$  (the ratio  $k_{43}/k_{34}$ ) is  $0.1 \text{ M}^2$  for wt SGLT1 and  $0.07 \text{ M}^2$  for the mutants and is largely responsible for the  $\text{Na}^+$  affinity measured at low  $\alpha\text{MG}$  concentration. This ratio also has a large influence on the position of the  $V_{1/2}$  of the  $Q$ - $V$  curve. The  $K_{34}$  and the  $k_{23}$  values were modified for simulation of transferred charges through the mutant proteins to account for the faster  $\tau_{\text{slow}}$  at negative  $V_m$  and the new  $V_{1/2}$  of  $-30$  mV (instead of  $-50$  mV), which was observed for both mutants (26). These two simple modifications yielded the fit for the  $Q$ - $V$  curve of Fig. 6 B (left panel), shown for mutant C511A, and the  $\tau_{\text{slow}}$ - $V$  curve shown in Fig. 6 B (right panel) at 5 mM  $\alpha\text{MG}$ .

### Simulation of the effect of $\alpha\text{MG}$ on the pre-steady-state currents

To interpret the effects of  $\alpha\text{MG}$  on the pre-steady-state currents, appropriate values for the parameters  $k_{45}$ ,  $k_{54}$ , and  $k_{51}$  have to be determined. Our strategy was to start by establishing parameters that could explain the behavior of the transferred charges of wt SGLT1 in the presence of different  $\alpha\text{MG}$  concentrations. Our first criterion was that the transferred charge had to disappear in the presence of a saturating  $\alpha\text{MG}$  concentration. The  $k_{ij}$  of the four-state model established in the absence of  $\alpha\text{MG}$  were maintained constant, and we investigated the effects of the three new  $k_{ij}$  on the simulated  $Q$ - $V$  curve. We started with the parameters proposed by Loo et al. (11) (see Table 1) but needed to increase the value of  $k_{51}$  by up to 40-fold over the original value ( $2000$  vs.  $50 \text{ s}^{-1}$ ) to reproduce the charge disappearance observed at high  $\alpha\text{MG}$  concentration. It was also clear that the ratio  $k_{45}/k_{54}$  influenced the  $V_{1/2}$  of the  $Q$ - $V$  curve: an increase in this ratio shifts the  $V_{1/2}$  toward more positive  $V_m$ . However, the absolute values of  $k_{45}$  and  $k_{54}$  (and not only their ratio) were also important in the global behavior of the  $Q$ - $V$  curve as a function of  $\alpha\text{MG}$  concentration. With respect to the values proposed by Loo et al. (11), the values of  $k_{45}$  and  $k_{54}$  had to be increased by factors of 30 and 50, respectively.

For the mutants, we first established the values of the parameters  $k_{210}$ ,  $k_{23}$ , and  $K_{340}$  to account for the faster time constants and the positive shift in  $V_{1/2}$  in the absence of  $\alpha\text{MG}$ . Changes in  $k_{210}$  and  $k_{23}$  were necessary along with more modest

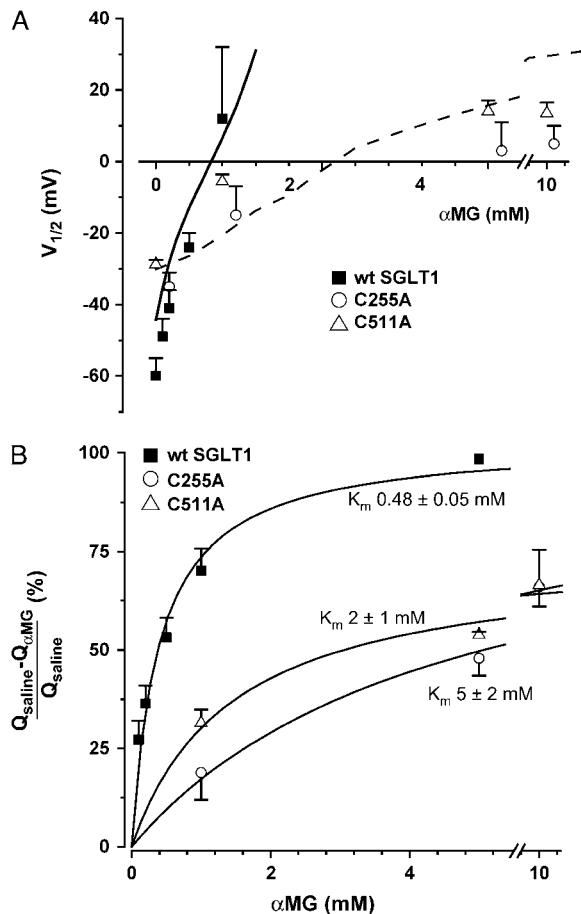


FIGURE 4 Effect of  $\alpha$ MG on  $V_{1/2}$  and estimation of  $K_m^{\alpha MG}$  with the transferred charge. (A)  $V_{1/2}$  of C255A (open circles), C511A (open triangles), and wt SGLT1 (solid squares) in the presence or absence of different  $\alpha$ MG concentrations. The lines represent the extracted  $V_{1/2}$  values from the theoretical Q-V curves obtained with the kinetic model's current simulations (solid line; wt SGLT1, dotted line: mutants). The phenomenological parameter  $V_{1/2}$  is obtained by a simple Boltzmann relation fitted to the experimental and theoretical Q-V curves (see Figs. 1 B, 2 B, and 3 B). (B) Normalized  $\alpha$ MG-dependent transferred charge in the presence of different  $\alpha$ MG concentrations ( $(Q_{\text{saline}} - Q_{\alpha\text{MG}})/Q_{\text{saline}}$ ) of wt SGLT1 (solid square), C255A (open circle), and C511A (open triangle) at  $-170$  mV. The line represents a simple Michaelis-Menten equation fitted to the points, and the corresponding  $K_m$  value from the fit is noted. Means  $\pm$  SE are shown.

changes in the remaining parameters describing the  $\alpha$ MG-independent steps of the kinetic model. The new values of  $k_{45}$ ,  $k_{54}$ , and  $k_{51}$  found for the wt transporter could not reproduce the observed Q-V and  $\tau_{\text{slow}}$ -V curves of the mutants. We found that the reactions describing  $\alpha$ MG binding and debinding had to be reduced by an order of magnitude and the reorientation of the fully loaded carrier ( $k_{51}$ ) had to be massively decreased from 2000 to 80  $\text{s}^{-1}$ . The best parameter set found is presented in Table 1 for wt SGLT1 and the mutants.

### Simulations of wt SGLT1

The simulated wt SGLT1 Q-V curves are shown in Fig. 6 A (middle panel), superimposed on the experimental data

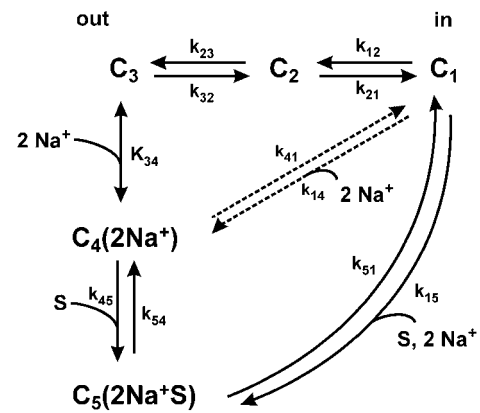


FIGURE 5 Kinetic model of the SGLT1 for the estimation of pre-steady-state currents. The voltage-independent substrate-binding/debinding events are included ( $k_{45}$ ,  $k_{54}$ ) in contrast to the models previously described in Chen et al. (4) and in Gagnon et al. (26). See Table 1 for rate constant values and Results and Discussion for further details.

points (in gray). Although, the shapes of the theoretical and experimental Q-V curves differ slightly, the general decrease due to external  $\alpha$ MG concentration is well represented. Three parameters were used to measure the accuracy of the model predictions: the  $V_{1/2}$  of the Q-V curves, the normalized transferred charge curves in the presence of different  $\alpha$ MG concentrations ( $(Q_{\text{saline}} - Q_{\alpha\text{MG}})/Q_{\text{saline}}$ ), as well as the  $\tau_{\text{slow}}$ -V curve. Fig. 4 A illustrates the phenomenological parameter  $V_{1/2}$  as a function of  $\alpha$ MG concentration extracted from the fits of a simple Boltzmann relation to our theoretical Q-V curves (solid line). Although the theoretical Q-V curve differs slightly from a simple Boltzmann relation, the model correctly represents the voltage shift produced by  $\alpha$ MG addition. The model also correctly predicts the total inhibition of the transferred charge by a saturating  $\alpha$ MG concentration. We estimated an apparent affinity for  $\alpha$ MG ( $K_m^{\alpha MG}$ ), using the remaining charge in the presence of  $\alpha$ MG ( $(Q_{\text{saline}} - Q_{\alpha\text{MG}})/Q_{\text{saline}}$ ), to compare it with that obtained experimentally. The parameters given in Table 1 yield a  $K_m^{\alpha MG}$  value of  $0.40 \pm 0.07$  mM at  $-170$  mV, which is not significantly different from the experimental value reported above. Finally, Fig. 6 A (right panel) shows that the  $\tau_{\text{slow}}$ -V curve values are close to the experimental values as the model reproduces very well the acceleration of the transient currents at hyperpolarizing  $V_m$  and shows the bell-shaped curve peak shifting toward more positive  $V_m$  as the  $\alpha$ MG concentration increases.

### Simulations of the mutant SGLT1s

The simulated Q-V curves for the mutant SGLT1s are shown in Fig. 6 B (middle panel), superimposed on the experimental data points for mutant C511A alone (in gray) because mutant C255A produced very similar values. The charge plateau value reached at hyperpolarizing  $V_m$ , at 5 and 10 mM  $\alpha$ MG, is reproduced very well by the modeled Q-V curve. The dotted line on Fig. 4 A illustrates  $V_{1/2}$  as a function of

**TABLE 1** Rate constants of a five-state kinetic model used for the pre-steady-state current simulations of wt SGLT1 and mutants

	$k_{120}^*$ (s <sup>-1</sup> )	$k_{210}$ (s <sup>-1</sup> )	$k_{23}$ (s <sup>-1</sup> )	$k_{320}$ (s <sup>-1</sup> )	$K_{340}^\dagger$ (M <sup>2</sup> )	$k_{41}^\ddagger$ (s <sup>-1</sup> )	$k_{45}$ (M <sup>-1</sup> s <sup>-1</sup> )	$k_{54}$ (s <sup>-1</sup> )	$k_{51}^\ddagger$ (s <sup>-1</sup> )
Loo et al. (11) <sup>§</sup>	25	600	—	—	0.021	0.3	10,000	20	50
Chen et al. (4)	790	370	100	280	0.153	—	—	—	—
wt SGLT1	790	370	110	130	0.1	10	300,000	1000	2000
Mutants	790	230	230	100	0.072	10	60,000	100	80

\*The constants  $k_{ij0}$  represent the value of  $k_{ij}$  at 0 mV (see Eq. 1). The rate constants that are independent of  $V_m$  are simply noted  $k_{ij}$ . The values of  $z_i$  used for the equivalent charge moving across the entire membrane electrical field in the step between state “ $i$ ” and “ $i + 1$ ” were  $-0.38$ ,  $-0.52$ , and  $-1.1$  for  $i$  varying from 1 to 3, respectively. The values for  $\alpha_i$  describing the asymmetry of the energy barrier were 0.3 and 0 for  $i$  of 1 and 2, respectively.

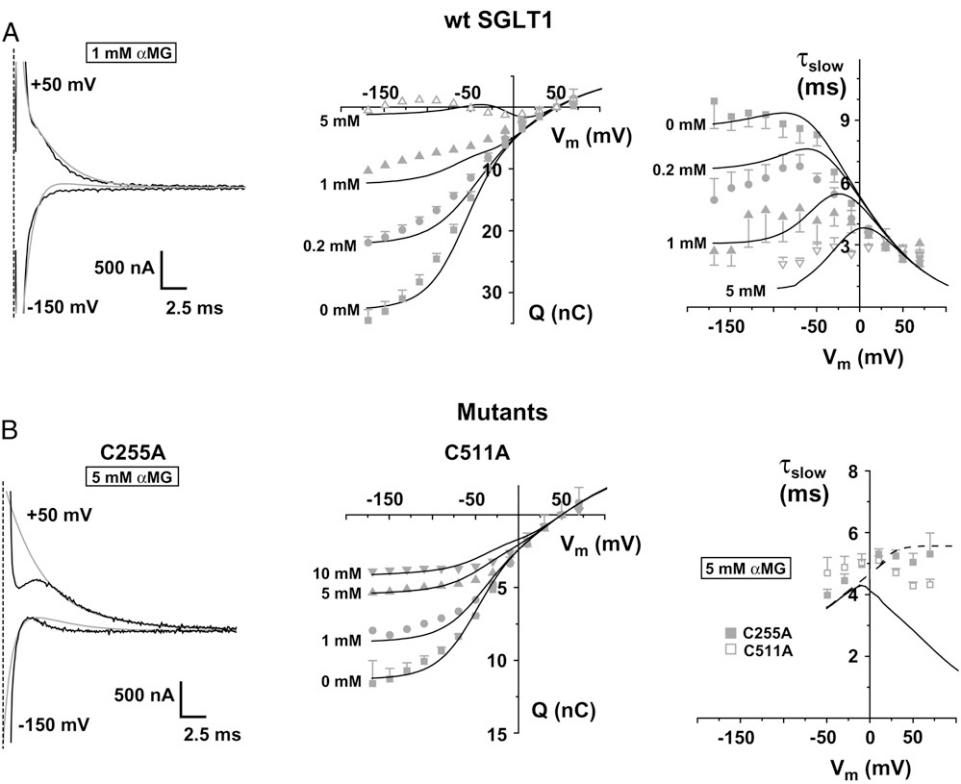
<sup>†</sup>The Na<sup>+</sup>-binding step is assumed to be in rapid equilibrium.  $K_{340}$  represents the ratio  $k_{430}/k_{340}$ , the equilibrium dissociation constant. This simplification eliminates the need for  $\alpha_3$  in the description of the voltage-dependent Na<sup>+</sup>-binding step.

<sup>‡</sup>The constants  $k_{14}$  (in M<sup>-2</sup>s<sup>-1</sup>) and  $k_{15}$  (in M<sup>-3</sup>s<sup>-1</sup>) were calculated to respect microreversibility. If the intracellular Na<sup>+</sup> concentration is set at 7 mM, the reaction rates from state 1 to state 4 are 0.009 s<sup>-1</sup> and 0.05 s<sup>-1</sup> for wt SGLT1 and for the mutants, respectively. If the intracellular  $\alpha$ MG concentration is 0.1 mM, the reaction rates from state 1 to state 5 are 0.05 s<sup>-1</sup> and 0.03 s<sup>-1</sup> for wt SGLT1 and for the mutants, respectively.

<sup>§</sup>In the model of Loo et al. (2,11), there are six states and no intermediate state between our  $C_1$  and  $C_3$ . Moreover, there are two additional states representing the intracellular facing fully loaded transporter and Na<sup>+</sup>-loaded transporter, which are lumped into  $k_{41}$  and  $k_{51}$  in our model. See Results and Discussion for more details and references.

concentration for the mutants. It is clear that the estimated  $V_{1/2}$  for the modeled Q-V curves of the mutants closely reproduced the characteristics of both mutants. The model accounts for the partial inhibition of the transferred charge at high  $\alpha$ MG concentrations. In addition, the  $K_{mQ}^{\alpha MG}$  was estimated with the remaining charge in the presence of  $\alpha$ MG ( $(Q_{\text{saline}} - Q_{\alpha MG})/Q_{\text{saline}}$ ) and provided the value of  $4 \pm 2$  mM at  $-170$  mV, which is close to the experimental values reported above for C255A and, to a lesser extent, for C511A ( $5 \pm 2$  mM and  $2 \pm 1$  mM, respectively). Finally, the theoretical  $\tau_{\text{slow}}$  values were

superimposed on the experimental values for both mutants at 5 mM  $\alpha$ MG in Fig. 6B (right panel). The model predicts two exponentials with significant amplitudes with time constants in the range of 2–6 ms in the presence of  $\alpha$ MG. The first one is almost identical with that observed in the absence of  $\alpha$ MG. The second one is slower at depolarizing  $V_m$  where it reaches a plateau value of  $\sim 5.5$  ms. Experimentally, a single exponential with a time constant in the millisecond range could be detected. Given the limited speed of the voltage pulse, the typical noise level found in our current recording, and given



**FIGURE 6** Predictions of the kinetic model. (A) Predictions of the kinetic model superimposed on the experimental data for wt SGLT1. The left panel shows simulated currents (gray line) superimposed on the experimental current (black line) at 1 mM  $\alpha$ MG for the indicated  $V_m$ ; the vertical dotted line indicates the beginning of the voltage step. The middle panel illustrates the Q-V curves, and the right panel shows  $\tau_{\text{slow}}$  as a function of  $V_m$  at different  $\alpha$ MG concentrations. (B) Predictions of the kinetic model superimposed on the experimental data for mutants. The left panel shows the experimental current (black line) of mutant C255A superimposed on the simulated currents at 1 mM  $\alpha$ MG (gray line) for the indicated  $V_m$ ; the vertical dotted line indicates the beginning of the voltage step. The middle panel illustrates the Q-V curves for mutant C511A, and the right panel shows  $\tau_{\text{slow}}$  as a function of  $V_m$  for both mutants at 5 mM  $\alpha$ MG. The currents from wt SGLT1 and C255A come from the experiment presented in Figs. 3A and 1A, respectively. The experimental data points were represented in gray for comparison. For clarity, only data points

from mutant C511A are presented for the Q-V curve. The Q-V curve was shifted vertically such that  $Q = 0$  at  $+50$  mV. As the model predicts significant amplitudes for two exponential components, the two time constants (a fast one, solid line, and a slower one, dashed line) are presented.



the fact that the two predicted time constants are in the same order of magnitude, it is conceivable that our experimental time constant would correspond to some intermediate value between the predicted ones. Thus it is concluded that the model reproduces fairly well the experimental time constants measured in the presence of  $\alpha$ MG.

## DISCUSSION

The identification of a disulfide bridge between C255 and C511 constitutes an important step in our understanding of how the 14 transmembrane segments are located with respect to each other and in the eventual identification of the physical structure that serves as the “voltage sensor” in SGLT1 (26). The two mutants C255A and C511A were found to display further interesting features which confirmed the importance of this region of the cotransporter. In this study we report that, in contrast with wt SGLT1, these two mutants exhibit pre-steady-state currents in the presence of a saturating  $\alpha$ MG concentration. By analyzing the dose-dependent effects of  $\alpha$ MG on the pre-steady-state currents of these mutants as well as for wt SGLT1, we sought to identify a satisfying kinetic explanation for both the partial diminution of mutant pre-steady-state currents by  $\alpha$ MG and for the complete disappearance of the wt SGLT1 transient currents.

The pre-steady-state currents in the absence of  $\alpha$ MG have been studied using cut open oocytes exposed to various  $\text{Na}^+$  concentrations, and a simple four-state kinetic model (4) was found to be consistent with the amplitudes and the time constants ( $\tau_{\text{fast}}$  (<1 ms) and  $\tau_{\text{slow}}$  (1–10 ms)) of the experimentally determined pre-steady-state currents as a function of the external  $\text{Na}^+$  concentration. The presence of these two time constants was more recently confirmed for rabbit SGLT1 (7,8) and for hSGLT1 (10). In this last study, fluorescently labeled cotransporters were also used and a slower time constant of  $\sim 100$  ms was reported in addition to  $\tau_{\text{fast}}$  and  $\tau_{\text{slow}}$ . A seven-state model was suggested for the translocation of the free transporter and the binding of two external  $\text{Na}^+$  ions, but the authors could not find a parameter set that would be in quantitative agreement with their own observations. Considering the time resolution provided by the two-electrode voltage-clamp technique, we decided to use the four-state model proposed by Chen et al. (4) to explain the effects of disrupting the disulfide bridge C255-C511 on the  $V_{1/2}$  of the Q-V and  $\tau_{\text{slow}}$ -V curves in the absence of  $\alpha$ MG (26). In the original model, it was assumed that a single  $\text{Na}^+$  ion was involved in the pre-steady-state currents. To incorporate  $\alpha$ MG binding, and given that the cotransport stoichiometry is 2  $\text{Na}^+$ :1 glucose (31), we simply replaced the original rate constant for  $\text{Na}^+$  binding ( $k_{34}$ ) by  $k_{34}/[\text{Na}^+]$  to account for both  $\text{Na}^+$  ions and made it a second order rate constant in  $\text{M}^{-2}\text{s}^{-1}$ . As the extracellular  $\text{Na}^+$  concentration is constant in this study, further studies will have to test whether the model used is consistent with the effects of changes in external  $\text{Na}^+$  concentration.

## Occupancy probabilities in the presence and absence of $\alpha$ MG

In the absence of  $\alpha$ MG and at  $-50$  mV, the set of rate constants proposed in Table 1 leads to occupancy probabilities ( $C_i$ ) of 5%, 22%, 43%, and 30% (from  $i = 1$  to 4), indicating that 73% of the  $\text{Na}^+$ -binding sites are exposed outside either in a free or  $\text{Na}^+$ -bound state (4,26). Obviously this situation is highly voltage dependent and, at  $+70$  mV,  $C_1$  and  $C_2$  now represent 52% and 40% of the cotransporter conformations, respectively. If, from state  $C_4$  to state  $C_1$ , the total number of unitary charges that can move across the entire membrane electrical field is 2 ( $|z_1 + z_2 + z_3| = 2$ ), the occupancy probabilities at  $+70$  mV indicate that all but 11% of it has already moved into the inward facing configuration. Fig. 7 A presents the occupancy probabilities in the absence of  $\alpha$ MG for an extreme voltage step from  $+70$  to  $-150$  mV. Upon hyperpolarization,  $C_1$  rapidly transforms into  $C_2$  and the free binding sites exposed to the extracellular solution ( $C_3$ ) become  $\text{Na}^+$ -bound immediately. The step  $C_1 \rightarrow C_2$  is considered to be mainly responsible for the fast component to the transient current. In contrast, the following slow transformation of  $C_2$  into  $C_3$  (which is in equilibrium with  $C_4(2\text{Na}^+)$ ) is clearly responsible for the slow component of the observed transient currents. Fig. 7, B and C, presents the changes in occupancy probability for a similar voltage step but in the presence of  $\alpha$ MG at 1 or 5 mM. At  $+70$  mV, the starting probabilities are independent of the presence of  $\alpha$ MG in the extracellular solution as are the fast events occurring in the first millisecond after hyperpolarization. At  $-150$  mV, the slowest rate constant in the reactions leading to the inward  $\text{Na}^+$ /glucose current is clearly  $k_{23}$ . This is why  $C_2$  accumulates transiently (75%) then relaxes to a value consistent with the steady-state cotransport rate allowed by the external  $\alpha$ MG concentration. This is shown in Fig. 7 D where the probability of finding  $C_2$  is plotted as a function of time for different  $\alpha$ MG concentrations.

For the mutants, the rate constants of Table 1 lead to occupancy probabilities of 2%, 9%, 45%, and 44% (from  $i = 1$  to 4) at  $-50$  mV in the absence of  $\alpha$ MG, which is quite similar to the  $C_i$  probabilities found for wt SGLT1. At  $+70$  mV, the distribution is slightly different from the wt SGLT1 as  $C_2$  now dominates with a probability of 42% with respect to  $C_1$  (34%), whereas the outward facing, free binding site ( $C_3$ ) presents a significant probability of 23%. Fig. 8 A presents the changes in  $C_i$  as a function of time for an extreme voltage pulse from  $+70$  mV to  $-150$  mV. Once again, in the absence of  $\alpha$ MG,  $C_2$  is transiently accumulated before relaxing to <5% as the  $\text{Na}^+$ -bound form ( $C_4(2\text{Na}^+)$ ) progressively rises to more than 90%. Fig. 8, B and C, depicts the occupancy probabilities in the presence of 1.5 and 5 mM  $\alpha$ MG. Under these circumstances, and in marked contrast to wt SGLT1,  $C_2$  continues to relax to a low value and it is  $C_5(2\text{Na}^+\text{S})$ , the fully loaded transporter, that progressively increases and attains up to 48% (this value increases to 58% at 10 mM  $\alpha$ MG). As

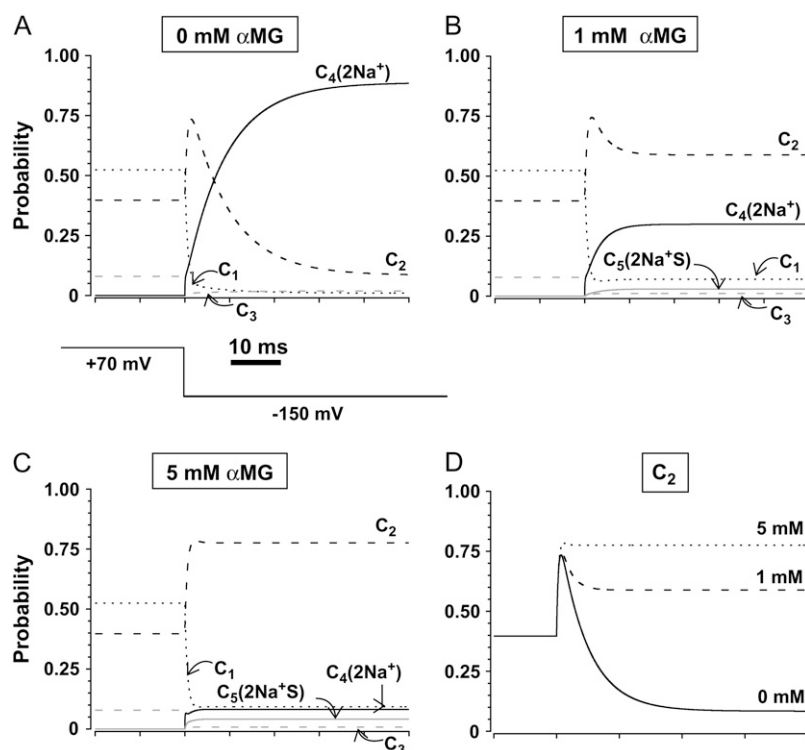


FIGURE 7 Occupancy probability ( $C_i$ ) as a function of time as calculated by the five-state kinetic model for wt SGLT1. (A) Time course of wt SGLT1 occupancy probabilities (90 mM  $\text{Na}^+$ ) for a  $V_m$  pulse from +70 mV to -150 mV in the absence of  $\alpha\text{MG}$  and in the presence of 1 mM  $\alpha\text{MG}$  (B) and 5 mM  $\alpha\text{MG}$  (C). (D) Time course of wt SGLT1  $C_2$  occupancy probability in the absence of  $\alpha\text{MG}$  (solid line, 90 mM  $\text{Na}^+$ ) or in the presence of 1 mM (dashed line) or 5 mM (small dashed line)  $\alpha\text{MG}$  for a  $V_m$  pulse from +70 mV to -150 mV.

illustrated in Fig. 8 D, contrary to what was seen for wt SGLT1, the  $C_2$  state increases (55%) and then relaxes to much lower steady-state values of 5% and 21% in the absence or presence of  $\alpha\text{MG}$ , respectively. This simply reflects the fact

that, for the mutant proteins, the slowest rate constant in the steps mediating  $\text{Na}^+$ /glucose cotransport at -150 mV is  $k_{51}$ . As the steps involved in generating the slow component of pre-steady-state currents are the transition between  $C_2$  and

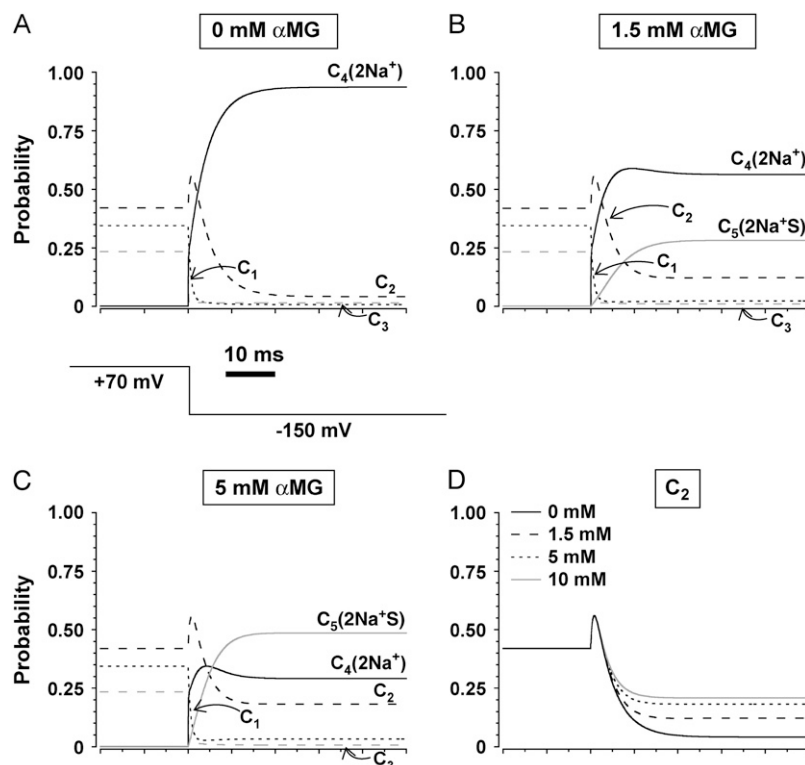


FIGURE 8 Occupancy probability ( $C_i$ ) as a function of time as calculated by the five-state kinetic model for the mutants. (A) Time course of mutant SGLT1 occupancy probabilities (90 mM  $\text{Na}^+$ ) for a  $V_m$  pulse from +70 mV to -150 mV in the absence of  $\alpha\text{MG}$ , in the presence of 1.5 mM  $\alpha\text{MG}$  (B), and in the presence of 5 mM  $\alpha\text{MG}$  (C). (D) Time course of mutant  $C_2$  occupancy probability in the absence of  $\alpha\text{MG}$  (solid line, 90 mM  $\text{Na}^+$ ) or in the presence of 1.5 mM (dashed black line) or 5 mM (small dashed black line) or 10 mM (solid gray line)  $\alpha\text{MG}$  for a  $V_m$  pulse from +70 mV to -150 mV.

$C_4(2Na^+)$ , Figs. 7 D and 8 D illustrate the reason transient currents disappear in the presence of  $\alpha$ MG for the wt transporter but not for the mutants.

### Apparent affinity for $\alpha$ MG

We presented two distinct methods of estimating the apparent affinity for the substrate: one can use the steady-state currents ( $I_{ss}(\alpha MG)$ ) and obtain  $K_{ml}^{\alpha MG}$  or the substrate-dependent charge disappearance ( $(Q_{saline} - Q_{\alpha MG})/Q_{saline}$ ) to obtain  $K_{mQ}^{\alpha MG}$ . For the wt SGLT1, both experimental (0.97 and 0.48 mM) and theoretical (0.36 and 0.40 mM) approaches show that the  $K_{ml}^{\alpha MG}$  and the  $K_{mQ}^{\alpha MG}$  are close in value. However, the two experimental  $K_m$  estimates for the mutants are significantly different, particularly for mutant C511A. It is important to specify that these two  $K_m$  are apparent  $K_m$  and depend not only on the rate constants  $k_{45}$  and  $k_{54}$  but also on the other rate constants. It seems that the rate-limiting step position is crucial for this discrepancy and that the two methods of obtaining an apparent affinity constant for  $\alpha$ MG should be considered with caution. The accordance between their values for wt SGLT1 may simply be coincidental.

### Role of the disulfide bridge C255-C511 in SGLT1

In a previous study, we have shown that the breakage of a disulfide bridge between C255 and C511 using dithiothreitol or by disruption through specific alanine mutations led to a displacement of the equilibrium position of the “voltage sensor” and to an acceleration of time constant of pre-steady-state current in the absence of  $\alpha$ MG (26). In this study, we established that the disulfide bridge C255-C511 (in hSGLT1) also plays a major role in facilitating the conformational change of the fully loaded cotransporter. In addition, a minor role was also detected in the  $\alpha$ MG-binding and -debinding reactions. In the absence of a tridimensional structure, it is impossible to know the exact position of this disulfide bridge in relation to the  $Na^+$  or  $\alpha$ MG-binding sites, but it is certainly important for the mechanical structure involved in those processes.

### CONCLUSIONS

In summary, this study has provided a quantitative explanation for the observation that transient currents disappear in the presence of  $\alpha$ MG for the wt SGLT1 but not for the mutant transporters. In wt SGLT1 and in the presence of substrate, the rate-limiting step is from state 2 to state 3. The transferred charges are not observed in this case because, upon hyperpolarization from a very positive to a very negative  $V_m$ , the large steady-state current requires a high  $C_2$  probability. Under these circumstances, the steady-state transporter distribution is predicted to simply move from state 1 to state 2, which should generate only a very fast transient current ( $\tau \approx 0.5$  ms). In contrast, for the mutants C255A and C511A in the presence of  $\alpha$ MG, the rate-limiting step is from state 5 to state

1. The transporter after having reached a high  $C_2$  probability will relax to a much lower level to reach the required  $C_5$  probability to account for the steady-state current. As the transporter moves from state 2 to states 3–5 through electrogenic steps, a slow transient current is generated. The behavior of the mutants underscores the role played by the rate-limiting step in the possibility of observing pre-steady-state currents. It also reveals the importance of the disulfide bridge C255-C511 in facilitating the translocation of the fully loaded transporter from the outward facing to the inward facing configuration.

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